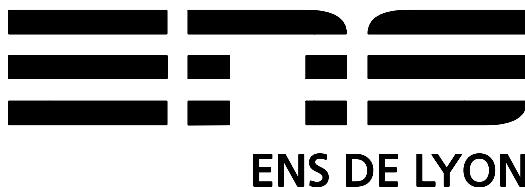


Quantum aspects of magnetic fragmentation in pyrochlore spin ice

Flavien Museur, Peter Holdsworth, Elsa Lhotel

VaQuM online school, July 2020



Quantum spin ice

- XXZ Hamiltonian on the pyrochlore lattice:

$$H = \sum_{\langle ij \rangle} J_{zz} S_i^z S_j^z - J_{\pm} (S_i^+ S_j^- + S_i^- S_j^+)$$

- Classical spin ice $J_{\pm} = 0$: frustration with local constraint

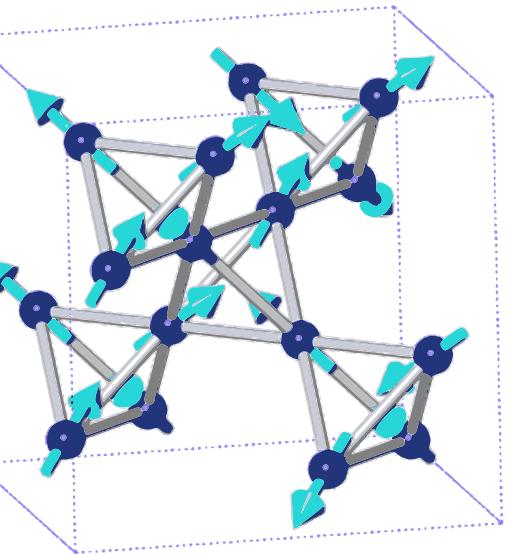
$$\mathcal{Q}_{\mathbf{r}} = \eta_{\mathbf{r}} \sum_{\mu} S_{\mathbf{r}\mu}^z = 0$$

- Mapping to a compact U(1) lattice gauge theory [2], [5] :

$\Phi_{\mathbf{r}}^\dagger, (\Phi_{\mathbf{r}})$ raising (lowering) operators for $\mathcal{Q}_{\mathbf{r}}$

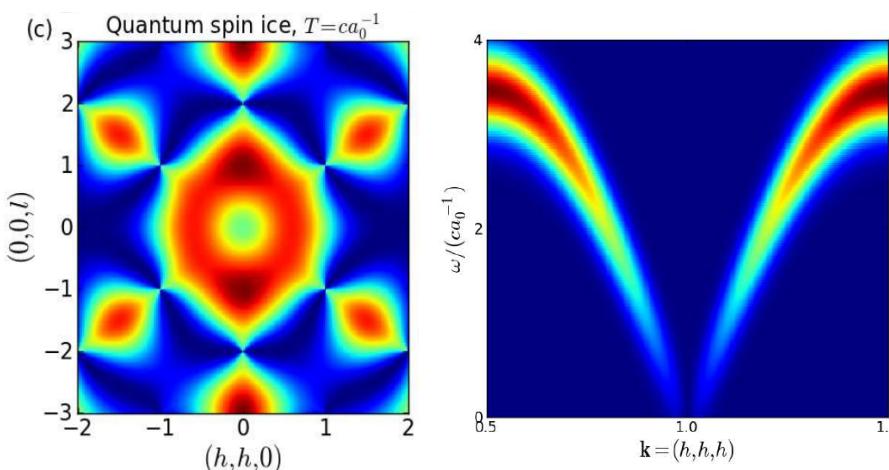
$$S_{\mathbf{r}\mu}^+ = \Phi_{\mathbf{r}}^\dagger e^{iA_{\mathbf{r}\mu}} \Phi_{\mathbf{r}+\hat{\mu}}$$

$$\Rightarrow H = \frac{J_{zz}}{2} \sum_{\mathbf{r}} \mathcal{Q}_{\mathbf{r}}^2 - J_{\pm} \sum_{\substack{\mathbf{r} \in A \\ \mu \neq \nu}} \Phi_{\mathbf{r}+\hat{\mu}}^\dagger \Phi_{\mathbf{r}+\hat{\nu}} e^{-i(A_{\mathbf{r}\mu} - A_{\mathbf{r}\nu})} - J_{\pm} \sum_{\substack{\mathbf{r} \in B \\ \mu \neq \nu}} \Phi_{\mathbf{r}-\hat{\mu}}^\dagger \Phi_{\mathbf{r}-\hat{\nu}} e^{i(A_{\mathbf{r}\mu} - A_{\mathbf{r}\nu})}$$



$i, j \in$ Pyrochlore lattice
 $\mathbf{r} \in$ Diamond lattice
 with sublattices A, B

On a bigger Hilbert space
 $\mathcal{H}_{tot} = \mathcal{H}_{\mathcal{Q}} \bigotimes \mathcal{H}_A$
 with Gauss' Law constraint



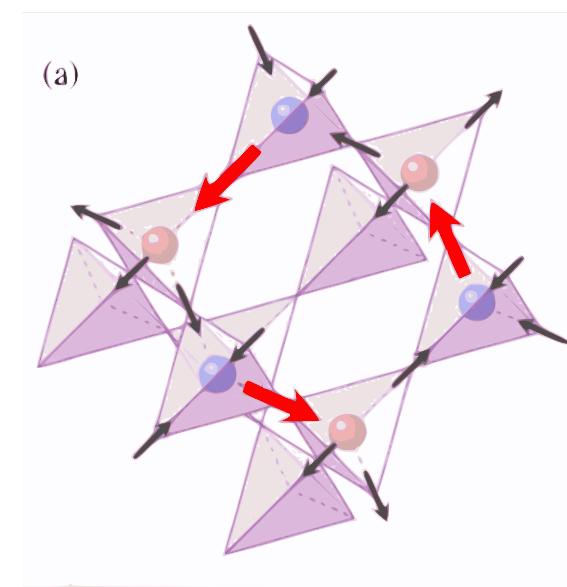
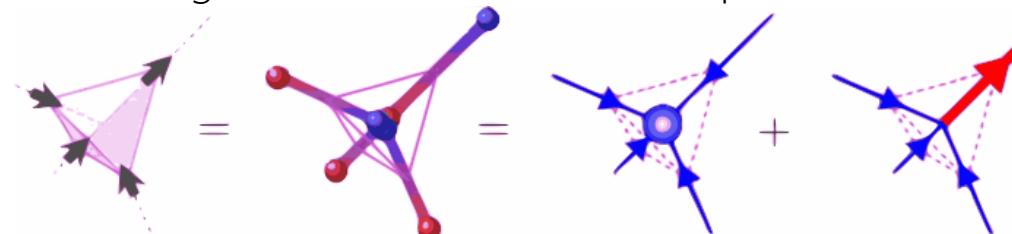
- Perturbative limit $J_{\pm} \ll J_{zz}$: spin loop terms \rightarrow U(1) spin liquid [11]
- Emergent EM with photon-like excitation
- Signatures in magnetic structure factor [4]

Fragmentation in spin ice [3]

- A staggered chemical potential can stabilize a monopole crystal:

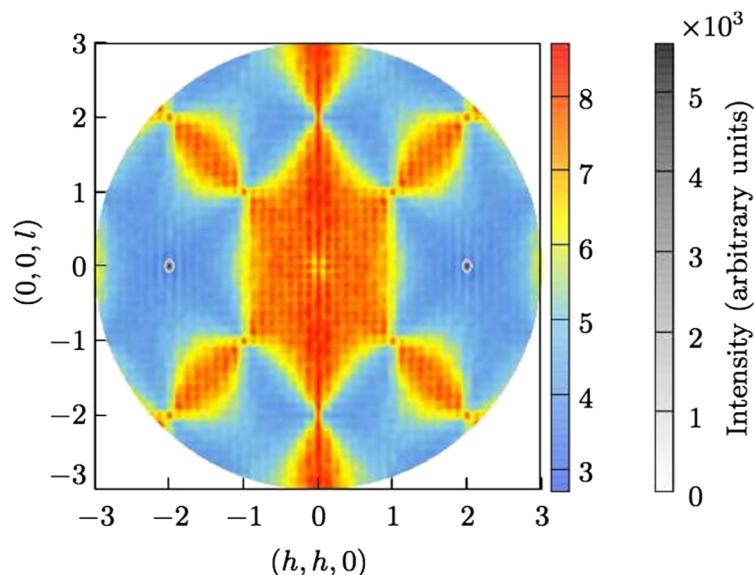
$$H = \frac{J_{zz}}{2} \sum_{\mathbf{r}} Q_{\mathbf{r}}^2 - J_{\pm} [\dots] - \Delta \sum_{\mathbf{r}} \eta_{\mathbf{r}} Q_{\mathbf{r}}$$

- Fragmentation of magnetization in a divergence full (longitudinal) and a divergence free (transverse) component



$$[M_{\mathbf{r}\mu}] \frac{a}{m} = (-1, -1, -1, 1) = \left(-\frac{1}{2}, -\frac{1}{2}, -\frac{1}{2}, -\frac{1}{2} \right) + \left(-\frac{1}{2}, -\frac{1}{2}, -\frac{1}{2}, \frac{3}{2} \right)$$

- Monopole order + emergent field with dipolar correlations (classically)
- One bond of strength $\pm 3/2$ per diamond site → mapping to a dimer model
- Extensive entropy ↔ closed loop of dimer moves



- Consequence on the classical structure factor: Braggs peak + diffuse pinch-points

Effective low-energy theories

➤ If we add quantum fluctuations, does the emergent dipolar field order ?

➔ Effective theory for the degenerate subspace of ground states

➤ Projection on the GS subspace satisfies the “Effective Schrödinger equation”:

$$\left[E_0 + \mathcal{P} H_1 \sum_{n=0}^{\infty} \left(\frac{1}{E - H_0} (1 - \mathcal{P}) H_1 \right)^n \mathcal{P} \right] |\Psi_0\rangle = H_{eff} |\Psi_0\rangle = E |\Psi_0\rangle$$

➤ Separation in diagonal and non-diagonal terms [7], [8] :

Diagonal terms

- Energy of a particular configuration of spins / dimers in the classically degenerate manifold
- At lowest order : potential energy of having a flippable 6-link loop $\mu \sim J_{\pm}^6/J_{zz}^5$

Non-diagonal terms

- Energy cost of tunnelling to another configuration
- At lowest order: kinetic energy associated with flipping a 6-link loop $g \sim J_{\pm}^3/J_{zz}^2$

➤ For quantum spin ice:

$$H_{eff,s} = \mu_s \sum_{\textcirclearrowleft} |\textcirclearrowleft\rangle\langle\textcirclearrowleft| + |\textcirclearrowright\rangle\langle\textcirclearrowright| - g_s \sum_{\textcirclearrowleft} |\textcirclearrowleft\rangle\langle\textcirclearrowright| + |\textcirclearrowright\rangle\langle\textcirclearrowleft|$$

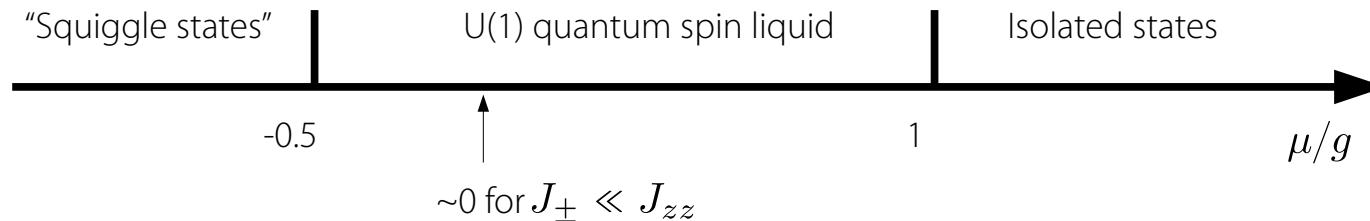
➤ For fragmented spin ice and emergent dimers:

$$H_{eff,d} = \mu_d \sum_{\textcirclearrowleft} |\texttriangledown\rangle\langle\texttriangledown| + |\texttriangle\rangle\langle\texttriangle| - g_d \sum_{\textcirclearrowleft} |\texttriangledown\rangle\langle\texttriangle| + |\texttriangle\rangle\langle\texttriangledown|$$

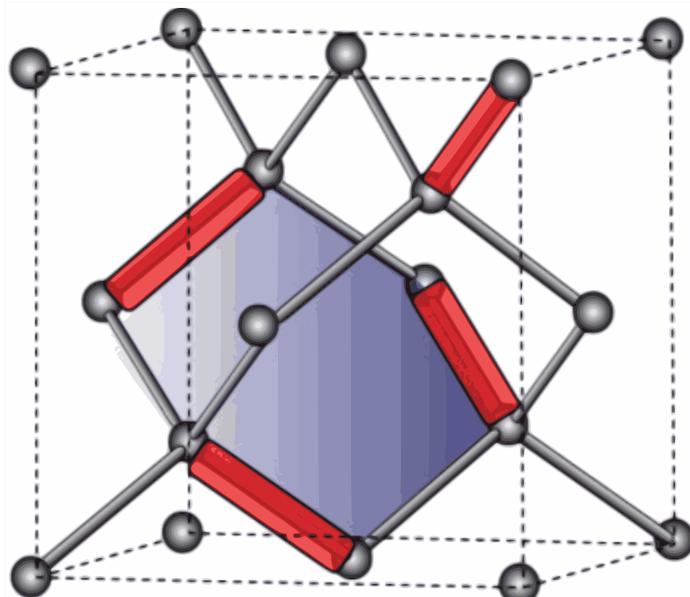
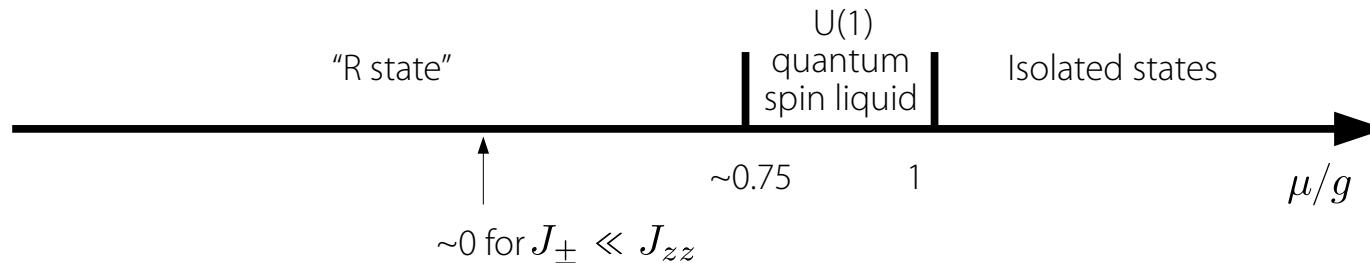
Possible quantum phases

➤ Depends on the value of $\mu/g \sim (J_{\pm}/J_{zz})^3$

➤ For quantum spin ice [4]:



➤ For quantum dimers [6]:



- In fragmented spin ice, **the emergent dipolar field is likely ordered**, with dimers maximising the number of flippable plaquettes
- Experimental signatures: structure factor of R state as one sector of the dimers + phase transition as $T \rightarrow 0$

Bibliography

- [1] T. A. Bojesen and S. Onoda, Physical Review Letters 119, (2017).
- [2] Z. Hao, A. G. R. Day, and M. J. P. Gingras, Physical Review B 90, (2014).
- [3] M. E. Brooks-Bartlett, S. T. Banks, L. D. C. Jaubert, A. Harman-Clarke, and P. C. W. Holdsworth, Physical Review X 4, (2014).
- [4] N. Shannon, O. Sikora, F. Pollmann, K. Penc, and P. Fulde, Physical Review Letters 108, (2012).
- [5] L. Savary and L. Balents, Physical Review Letters 108, (2012).
- [6] O. Sikora, N. Shannon, F. Pollmann, K. Penc, and P. Fulde, Physical Review B 84, (2011).
- [7] D. L. Bergman, R. Shindou, G. A. Fiete, and L. Balents, J. Phys.: Condens. Matter 19, 145204 (2007).
- [8] D. L. Bergman, R. Shindou, G. A. Fiete, and L. Balents, Phys. Rev. B 75, 094403 (2007).
- [9] R. Moessner, S. L. Sondhi, and M. O. Goerbig, Phys. Rev. B 73, 094430 (2006).
- [10] S. V. Isakov, K. Gregor, R. Moessner, and S. L. Sondhi, Physical Review Letters 93, (2004).
- [11] M. Hermele, M. P. A. Fisher, and L. Balents, Physical Review B 69, (2004).