

Probing many-body entanglement via inverse statistical methods

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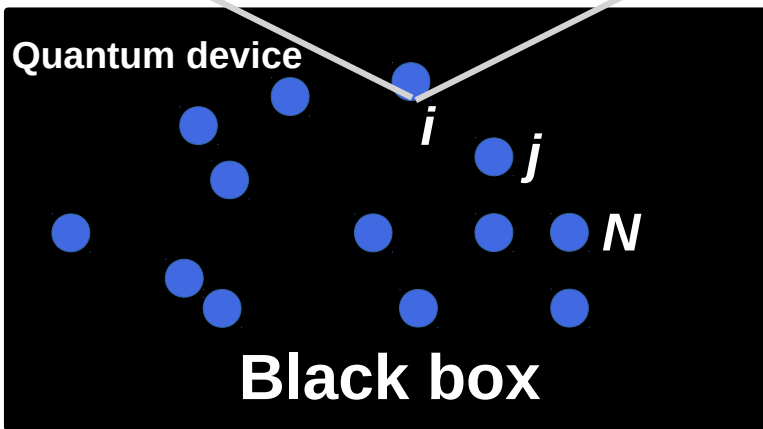
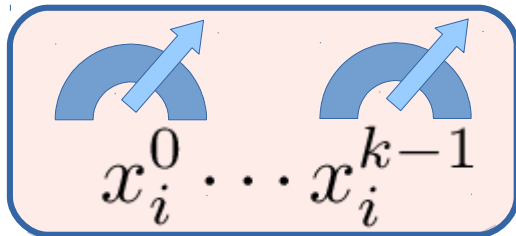
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Certifying entanglement via Bell's inequalities



Quantum data

$$\langle f(\sigma) \rangle_Q \quad (\sigma = \{x_i^a\})$$

(e.g. $\langle x_i^a \rangle_Q, \langle x_i^a x_j^b \rangle_Q, \dots$)

Bell's local-variable (LV) model (classical data):

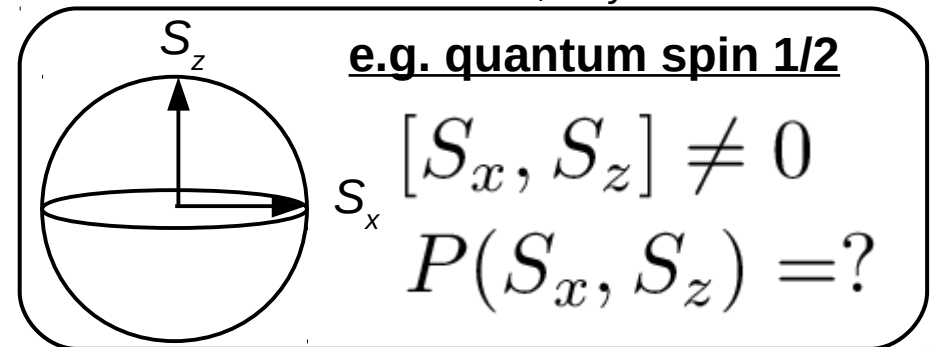
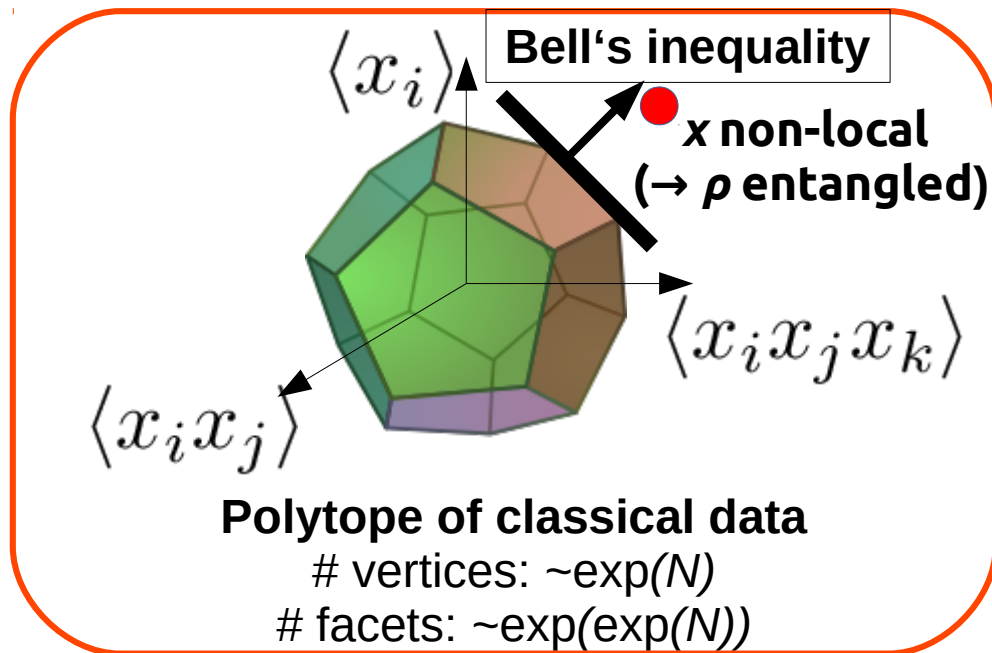
Global probability distribution $P_{LV}(\sigma)$ s.t.

$$\langle f_r(\sigma) \rangle_{LV} = \langle f_r(\sigma) \rangle_Q \quad (r = 1, \dots, R)$$

Always exists if the state is *not* entangled

J. S. Bell, Physics 1964

A. Fine, Phys. Rev. Lett. 1982



Membership problem:

NP hard

\rightarrow No scalable approach solving every instance

Statistical-inference approach to Bell's inequalities

IF & T. Roscilde, arxiv: 2004.07796 (2020)

State-of-the-art approaches:

- Brute force (linear program, solve convex hull problem → limited to small scenarios)
- Try some existing Bell inequalities (not generic)
- Remove structure from the data (symmetrize them; misses some data points)
Tura *et al* (Science 2014), Wang *et al* (PRL 2017), Fadel & Tura (PRL 2017)
- Approximate the polytope from the outside (hierarchy of semi-definite programs)
Baccari *et al* (PRX 2017)

Inverse problem: max entropy solution

$$P_{LV}(\sigma) = Z^{-1} \exp\left[\sum_{r=1}^R K_r f_r(\sigma)\right]$$

Boltzmann distribution over a classical model

E.T. Jaynes, Phys. Rev. 1957

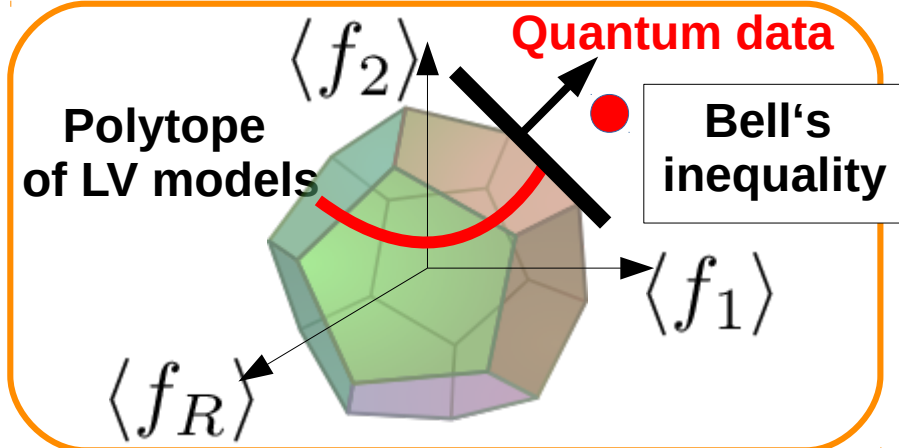
Update rule: iterate

$$K'_r = K_r - \epsilon(\langle f_r \rangle_{LV} - \langle f_r \rangle_Q)$$

gradient-descent of the log-likelihood (**convex**)

Nguyen *et al.*, Adv. Phys. 2017

**Trajectory of classical data
converging to the quantum data**

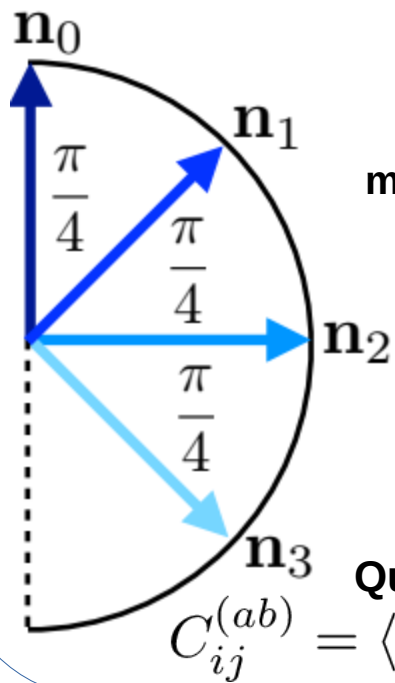


Generic: applies to any quantum data
Scalable: Monte-Carlo sampling of classical statistical models (unless e.g. spin glass)

Application: Heisenberg antiferromagnets

$$\mathcal{H} = J \sum_{\langle i,j \rangle} \vec{S}_i \cdot \vec{S}_j$$

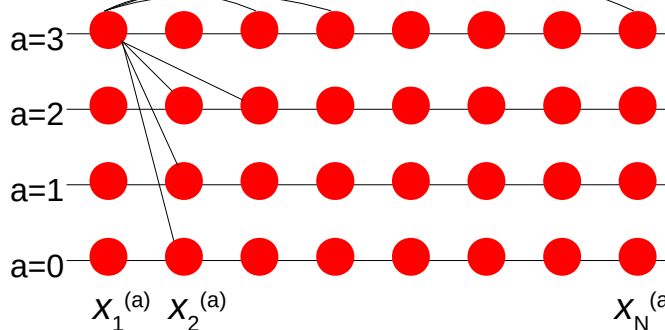
S=1/2 quantum spins on a square lattice
 Ground state: many-body singlet $S_{\text{tot}}=0$



Uniform measurement strategy (k=4)

Quantum data:
 $C_{ij}^{(ab)} = \langle \hat{\sigma}_i^{\mathbf{n}_a} \hat{\sigma}_j^{\mathbf{n}_b} \rangle_Q$

Fitting the quantum data with LV models



Equilibrium Ising model
 $H = - \sum_{i < j} \sum_{a,b} K_{ij}^{(ab)} x_i^{(a)} x_j^{(b)}$

Novel Bell inequality (permutationally invariant)

$$\langle B \rangle_{LV} = S_{11} + S_{12} + S_{22} + S_{23} + \dots + S_{kk} - S_{k1} \geq -2N(k-1)$$

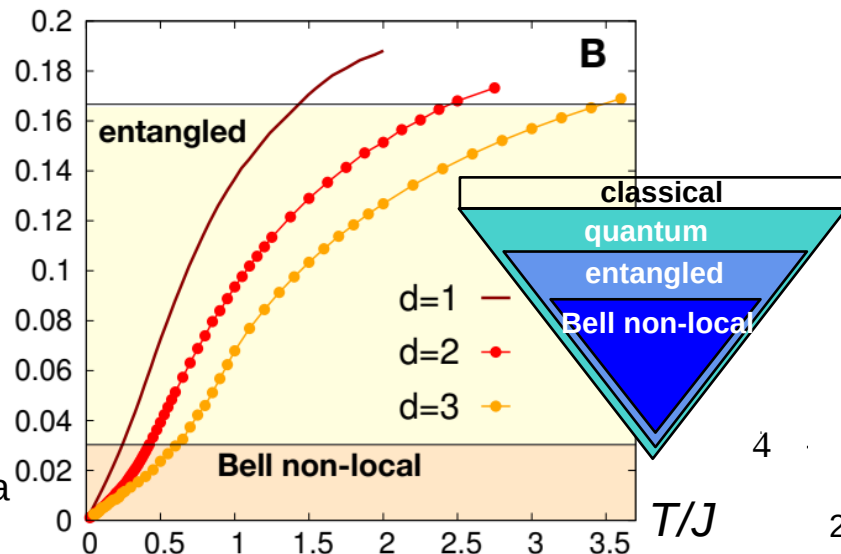
$$S_{ab} = \sum_{i \neq j} C_{ij}^{(ab)}$$

Many-body singlets: $\langle B \rangle_Q = -Nk[1 + \cos(\pi/k)]$
 (maximal quantum violation)

Assuming SU(2) invariance of the quantum state:

$$\langle \hat{J}_z^2 \rangle / N < \frac{1}{4} - \frac{k-1}{2k[1 + \cos(\pi/k)]} \quad (\text{Bell non-locality})$$

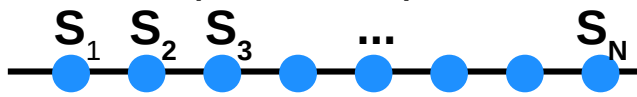
Compared to: $\langle \hat{J}_z^2 \rangle / N < \frac{1}{6}$ (entanglement, Tóth et al. PRA 2004)



Similar structure to Schmied et al (Science 2016); emerges from the data
 IF & T. Roscilde, arxiv: 2004.07796 (2020)

Extension: solving to the separability problem

N quantum spin-1/2

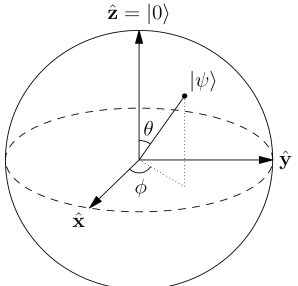


$$\rho_{\text{sep}} = \prod_{i=1}^N \int_{|\mathbf{S}_i|=1} d\mathbf{S}_i P(\mathbf{S}_1, \dots, \mathbf{S}_N) \rho_1(\mathbf{S}_1) \otimes \dots \otimes \rho_N(\mathbf{S}_N)$$

with $\rho(\mathbf{S}_i) := |\psi(\mathbf{S}_i)\rangle\langle\psi(\mathbf{S}_i)|$

P : statistical distribution over classical rotators

Bloch sphere



$$|\mathbf{S}(\theta, \phi)|^2 = 1$$

$$|\psi(\mathbf{S})\rangle = \cos(\theta/2)|\uparrow\rangle + \sin(\theta/2)e^{i\phi}|\downarrow\rangle$$

Quantum data = 2-body correlation matrix:

$$C_{ij}^{ab} := \langle \hat{\sigma}_i^a \hat{\sigma}_j^b \rangle \equiv_{(\rho_{\text{sep}})} \langle S_i^a S_j^b \rangle_P$$

$$P(\{\mathbf{S}_i\}) \propto \exp\left[\sum_{i,j=1}^N \sum_{a,b \in \{x,y,z\}} K_{ij}^{ab} S_i^a S_j^b \right]$$

Solution: Generalized XYZ model for classical rotators

Tomographically complete
(no other possible entanglement witness based on 2-body correlators)

General case: local Hilbert space dim = d:
- 2d - 2 classical variables
- d² - 1 measurements (e.g. generators of SU(d))

(unpublished)